

***The Broadening of Spectral Lines
by Collisions with Neutral
Hydrogen Atoms in Cool Stars***

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Scientific Motivation

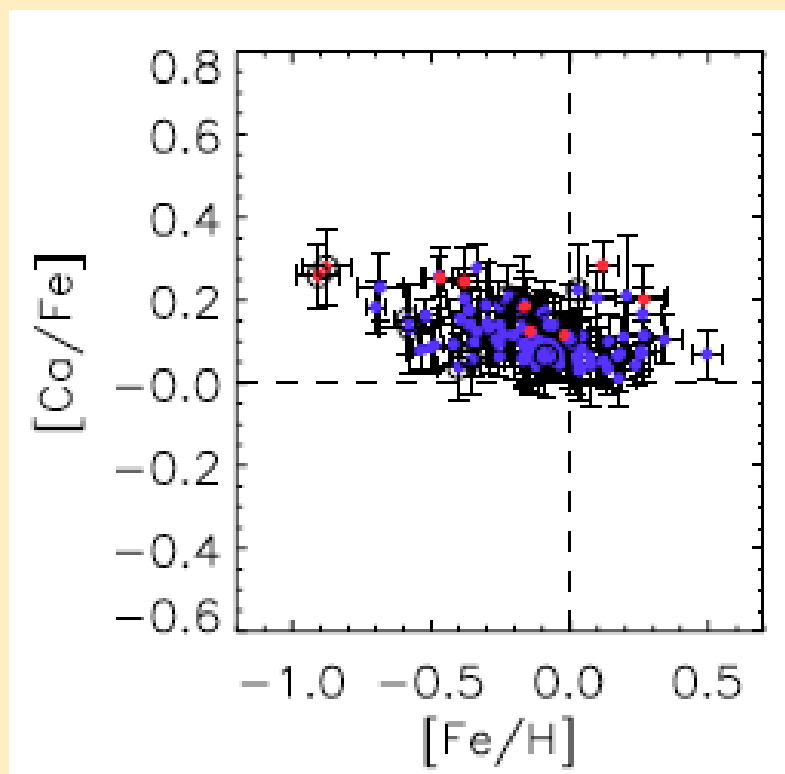
Need accurate stellar chemical abundances from spectroscopy for many problems in modern astrophysics:

- Understanding the chemical and dynamical evolution of the Galaxy, e.g.:
 - Origin and evolution of bulge, thin and thick disks
 - Merger history
 - Astrophysical sites of nucleosynthesis processes
- Solar composition and its place in the solar neighbourhood
- Strong lines are often the best to use - in some cases, e.g. very cool stars, galaxies, there is no choice

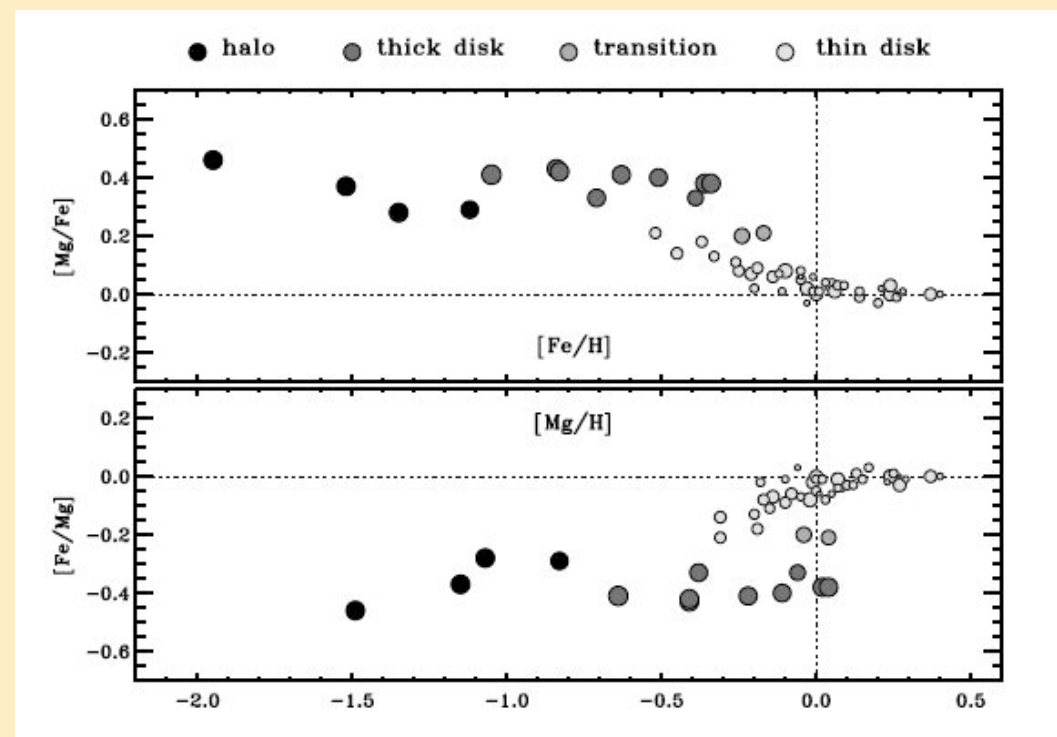


Scientific Motivation

Solar neighbourhood



Galactic Thick and Thin Disks



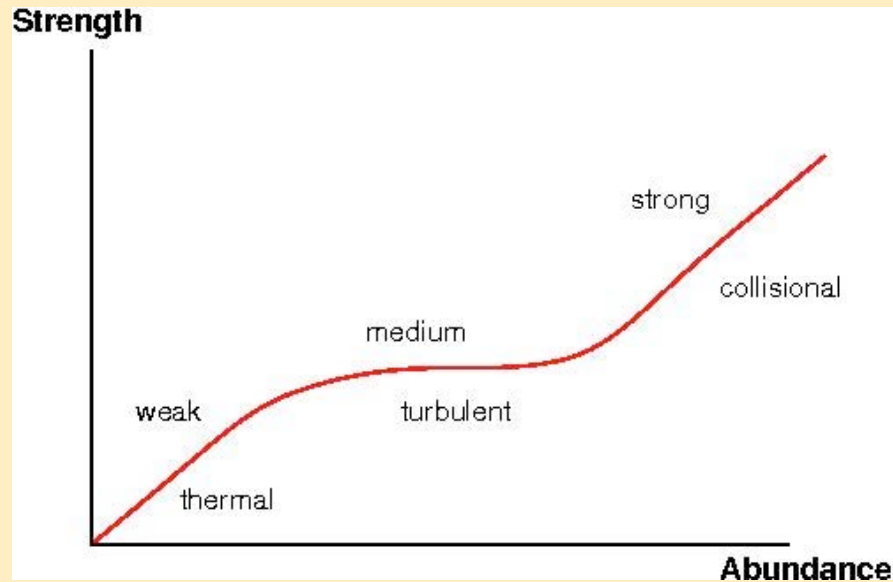
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Scientific Motivation



Weak lines are susceptible to blends

Medium strength lines are saturated and insensitive to abundance

Strong lines typically have the best oscillator strengths, and often no choice: E.g. very cool stars, galaxies, distant stars

Cool star atmospheres dominated by neutral hydrogen, in its ground state

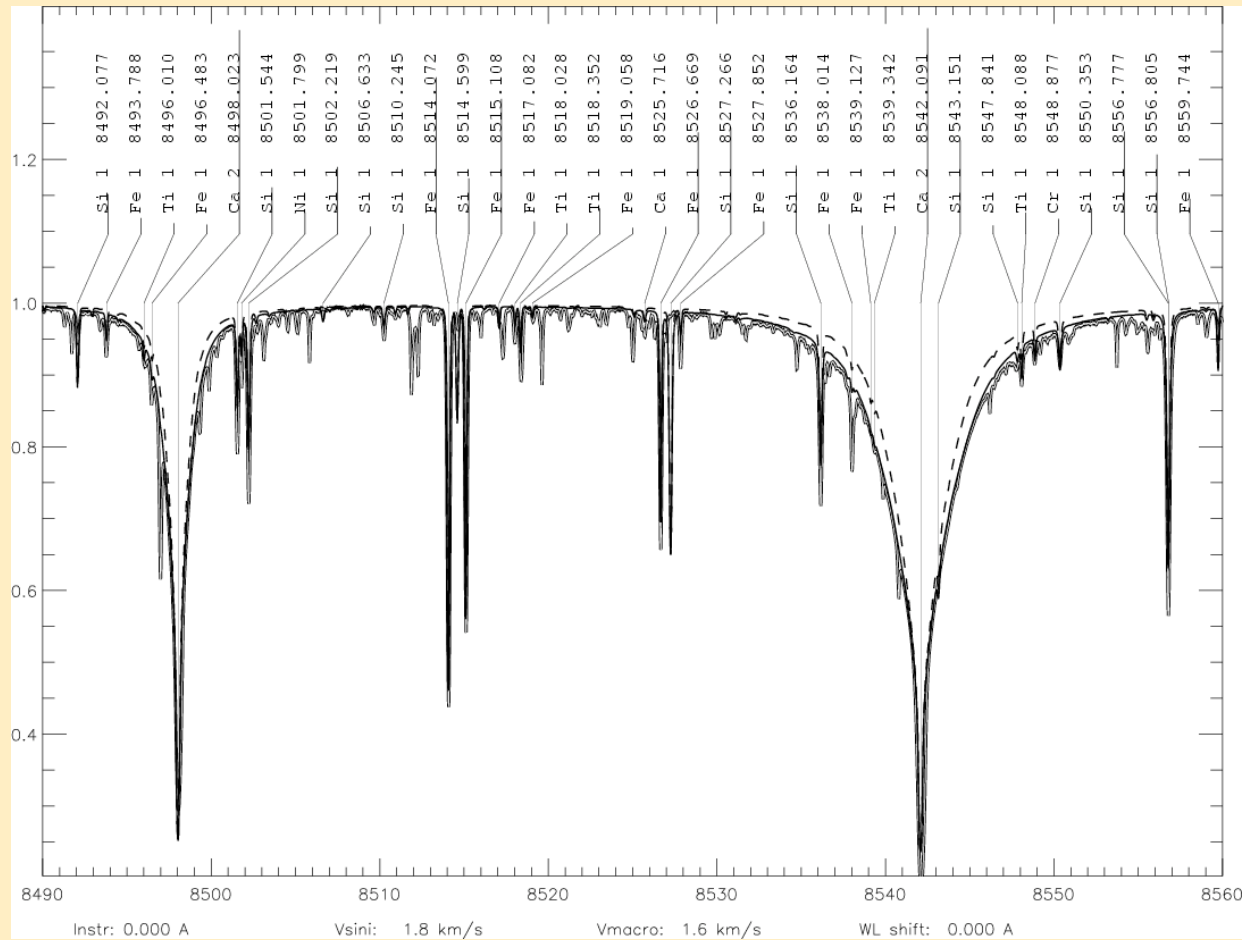
90% H, 9% He, 1% metals $\frac{N_H}{N_e} \approx 10^4$

Weight of number means generally dominant over electrons



Scientific Motivation

- E.g. Ca II IR triplet in the Sun



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History

- Lorentz (1906) Impulse theory (prior to QM)
- Weisskopf/London (1930's) VdW potential, strong collisions
- Lindholm/Foley (1940's) weak collisions, averaged interaction
- Unsöld (1950's) apply VdW + approx, general formula

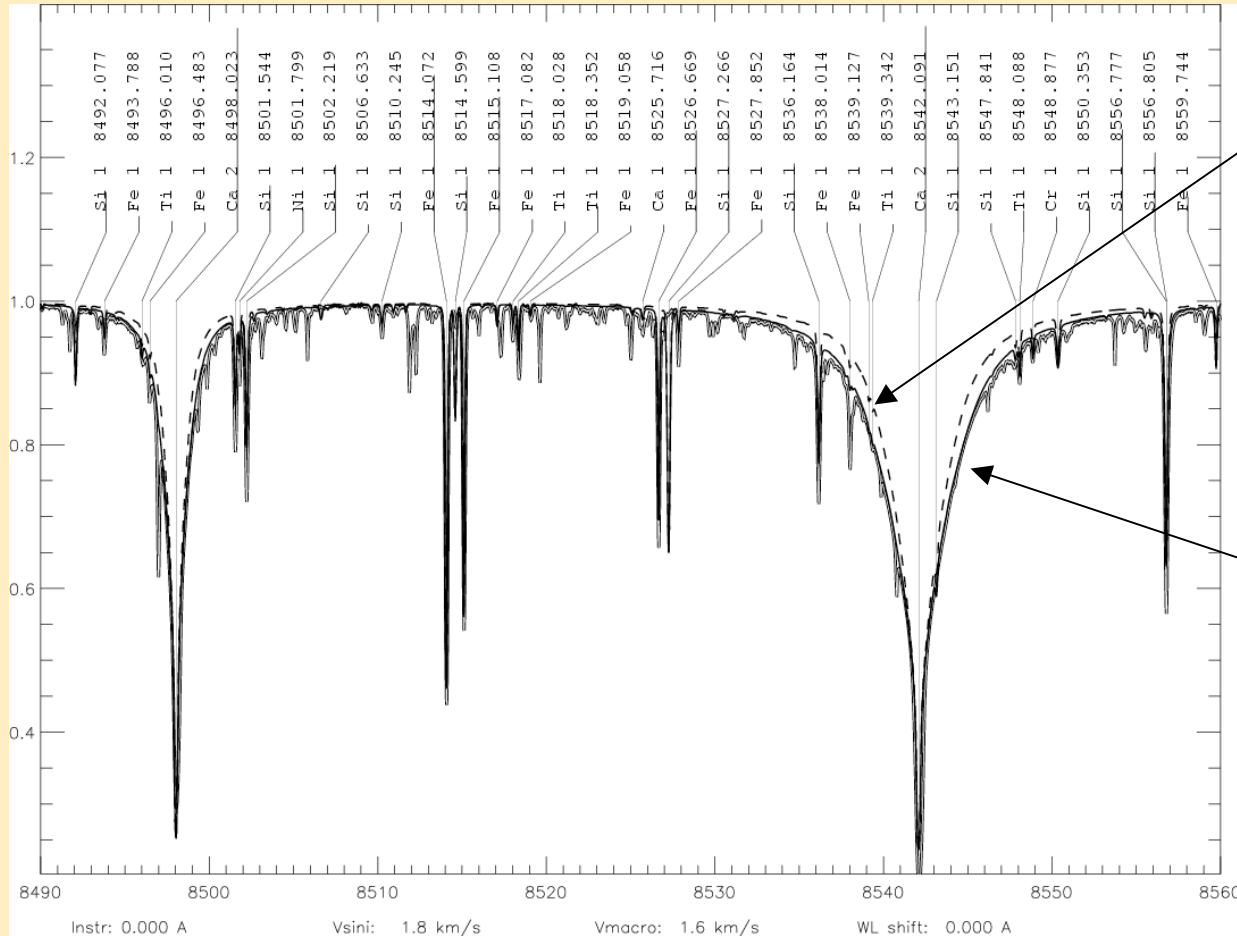
VdW = van der Waals $\Delta E \approx C_6 / R^6$

$$\Gamma = 17v^{3/5} C_6^{2/5} N_H$$

- (1960's -) Accurate calculations for a few specific cases



Scientific Motivation



Unsöld theory

Solar observations

Need a fudge factor about 1 to 4 to fit solar lines



Scientific Motivation

- Astrophysical evidence that Unsöld theory is inadequate
- Astrophysics needs a theory to compute data for a large number of lines of various elements (e.g. Fe, Ni, Mg, Ca, etc)
- Theory should ideally to be simple to use!



History

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- Weisskopf/London (1930's) VdW potential, strong collisions
- Lindholm/Foley (1940's) weak collisions, averaged interaction
- Unsöld (1950's) apply VdW + approx, general formula

- Brueckner (1970's) numerical RSU perturbation theory
- O'Mara (1970's) analytic RSU perturbation theory
- Anstee & O'Mara (1992) removal of averaged interaction+more
- Barklem & O'Mara (1998-) extension to d and f states, and ions

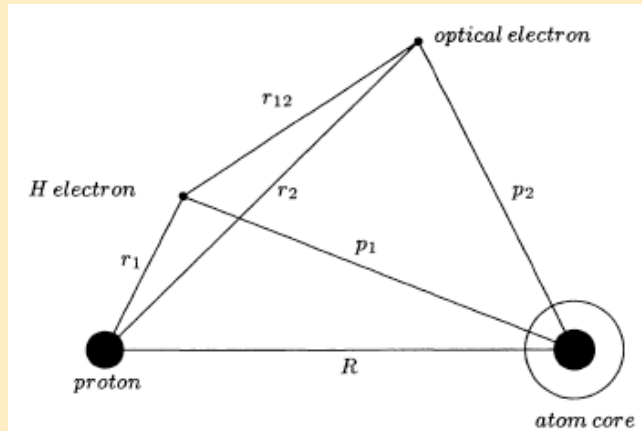


ABO theory

- isolated lines
- impact approximation (Lorentzian profile)
- classical straight path approximation
- no fine structure (orbital angular momentum basis, n/m)
- no quenching
- Important part is the potentials:



ABO theory: RSU potentials



$$V = \frac{1}{R} + \frac{1}{r_{12}} - \frac{1}{r_2} - \frac{1}{p_1}$$

Covalent partitioning
(Ionic partitioning neglected)

$$\Delta E = \langle i|V|i \rangle + \sum_{j \neq i} \frac{\langle i|V|j \rangle \langle j|V|i \rangle}{E_i - E_j}$$

Rayleigh-Schrödinger Perturbation Theory
to 2nd order (no exchange)

$$\approx \langle i|V|i \rangle + \frac{1}{E_p(R)} \langle i|V^2|i \rangle$$

← Unsöld approximations

$$\approx \langle i|V|i \rangle + \frac{1}{E_p} \int_0^\infty R_{nl}^2(p_2) I_{lm}(p_2, R) p_2^2 dp_2$$

$$E_p(R) \approx E_p(\infty)$$

Analytic functions where all integrals except the final integral
over the radial wavefunction have been performed



ABO theory: RSU potentials

$$\Delta E = \langle i|V|i \rangle + \sum_{j \neq i} \frac{\langle i|V|j \rangle \langle j|V|i \rangle}{E_i - E_j}$$

$$\approx \langle i|V|i \rangle + \frac{1}{E_p(R)} \langle i|V^2|i \rangle$$

$$\approx \langle i|V|i \rangle + \frac{1}{E_p} \int_0^\infty R_{nl}^2(p_2) I_{lm}(p_2, R) p_2^2 dp_2$$

- Coulomb wavefunctions (n^*, l)
- Unsöld approx $E_p = -2/\alpha_H = -4/9$
- Cross sections depend only on $n^*, l \longrightarrow$ Independent of species!

$$n^* = \left[2E_{binding} \right]^{-1/2}$$



ABO theory: Results

Cross section

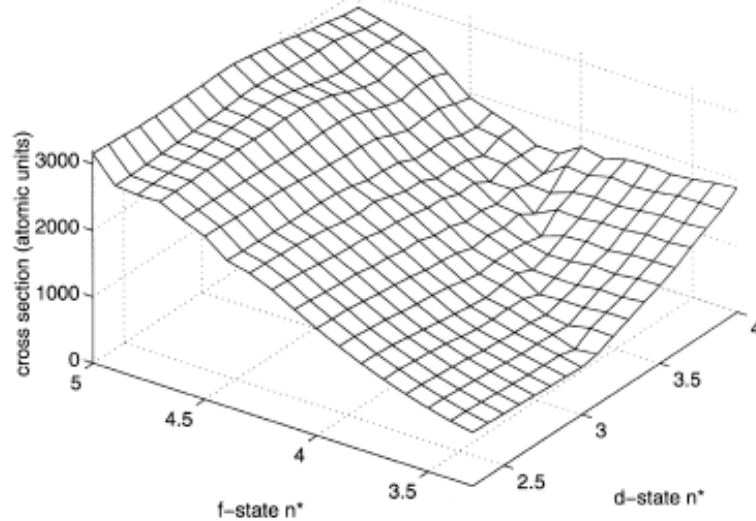


Figure 1. Plot of the cross-section for a perturber velocity of 10^4 m s^{-1} against the effective principal quantum number of the two states.

Velocity dependence

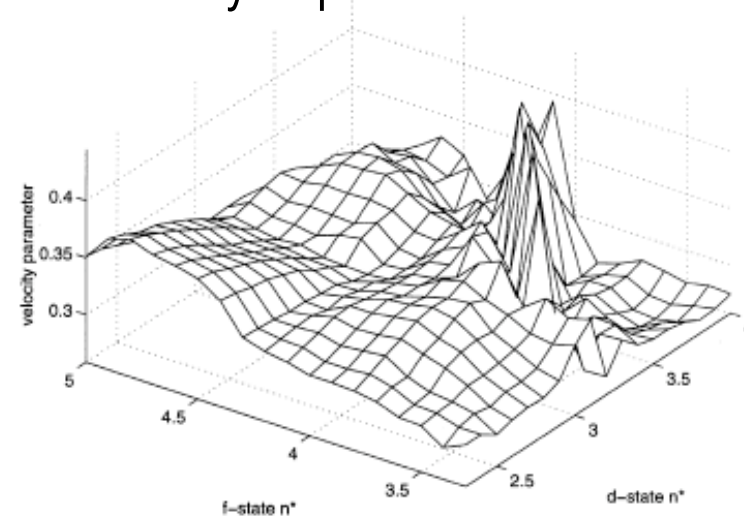


Figure 2. Plot of the velocity parameter α against the effective principal quantum number of the two states.

Tabulated results dependent only on effective principal quantum number

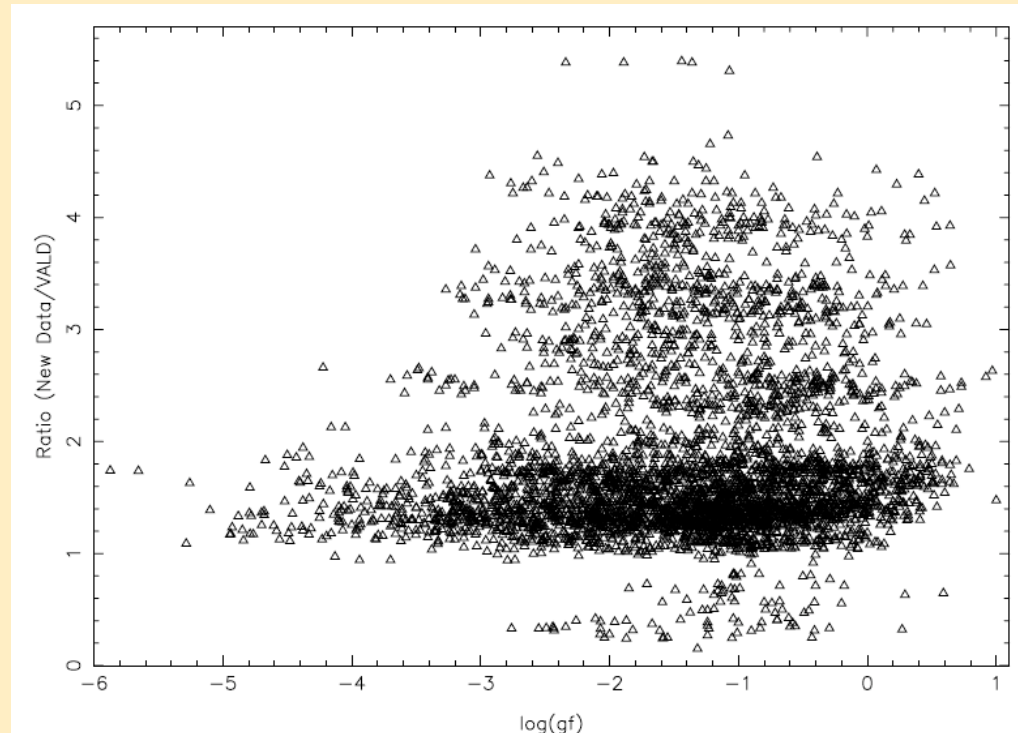
Simple velocity behaviour:

$$\sigma(v) = \sigma(v_0) \left(\frac{v}{v_0} \right)^{-\alpha}$$



ABO theory: Results

Comparison with Unsöld's theory for 4891 lines, Li to Ni:

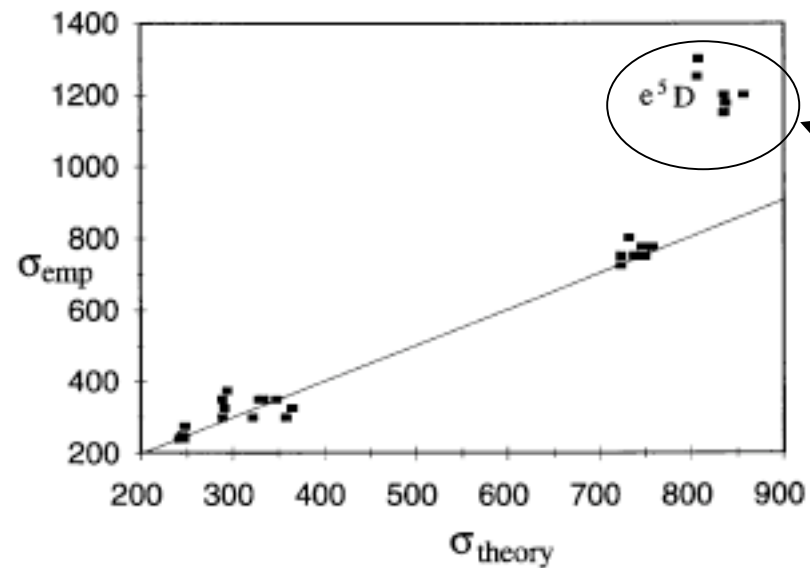


Mean difference of 1.88, in rough agreement with astrophysical fudge factor commonly used of 2



Applications

- Application to lines of Na, Ca and Fe in the solar spectrum indicate uncertainties as low as 5%



Strongly Stark
broadened

Figure 8. Empirical cross-sections obtained under the assumption that the solar abundance of iron is the same as in meteorites are plotted against theoretical cross-sections. The agreement is good except for lines with upper e^3D upper states.



Testing

- Comparison with more detailed calculations indicates uncertainties of order 5-20%.

Table 1. Comparison of line widths per unit perturber density (w/n_H in units $10^{-8} \text{ cm}^3 \text{ rad s}^{-1}$) at 5000 K for resonance lines of Mg, Ca and Sr using different potentials and dynamics.

Dynamics →	Quantal	Semi-classical	Semi-classical	Semi-classical
Potentials →	MOLPRO	MOLPRO	ABO	Hybrid
Mg	1.13	1.10	1.01	1.25
Ca	1.23	1.24	1.10	1.28
Sr	1.49	1.48	1.18	1.48

No difference due to dynamics

10-20% due to potentials
-neglect of ionic crossing



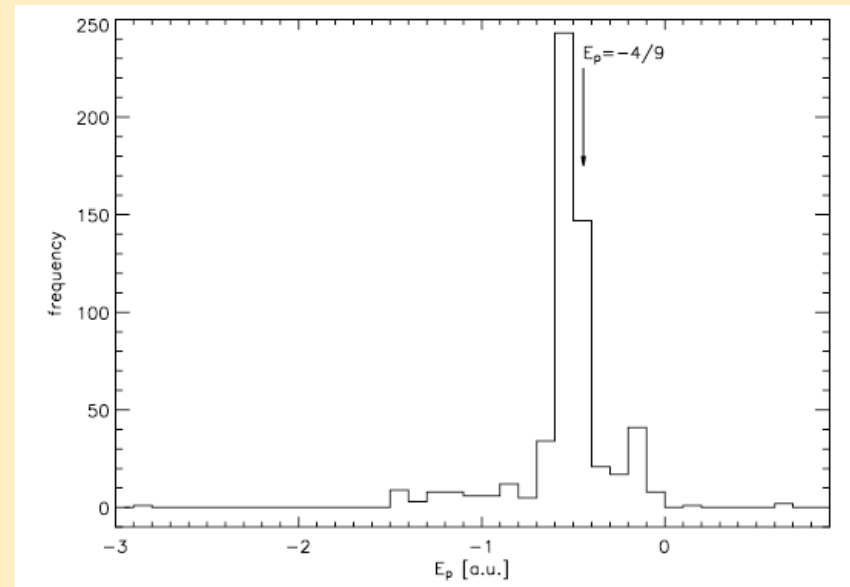
Extension to Ions

- Unsöld approx $E_p = -4/9$ is no longer valid
- Compute via

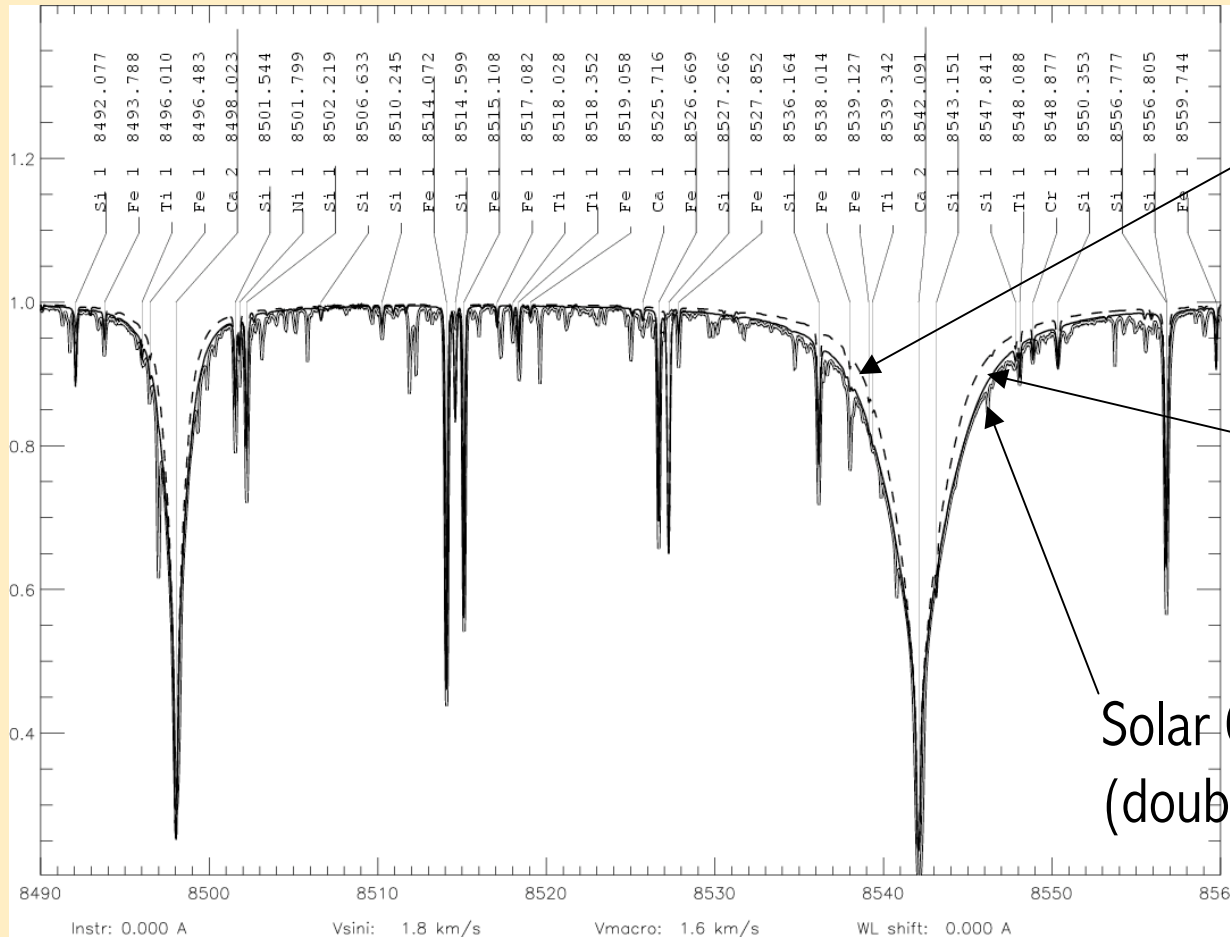
$$C_6 = \frac{3}{2} \sum_{k' \neq k} \sum_{l' \neq l} \frac{f_{kk'}^A f_{ll'}^H}{(\Delta E_{k'k}^A + \Delta E_{l'l}^H) \Delta E_{k'k}^A \Delta E_{l'l}^H},$$

$$E_p = -\frac{2 \langle p_2^2 \rangle}{C_6}.$$

Large Scale calculations for Fe II using large scale semi-empirical atomic data calculations by Kurucz



Applications



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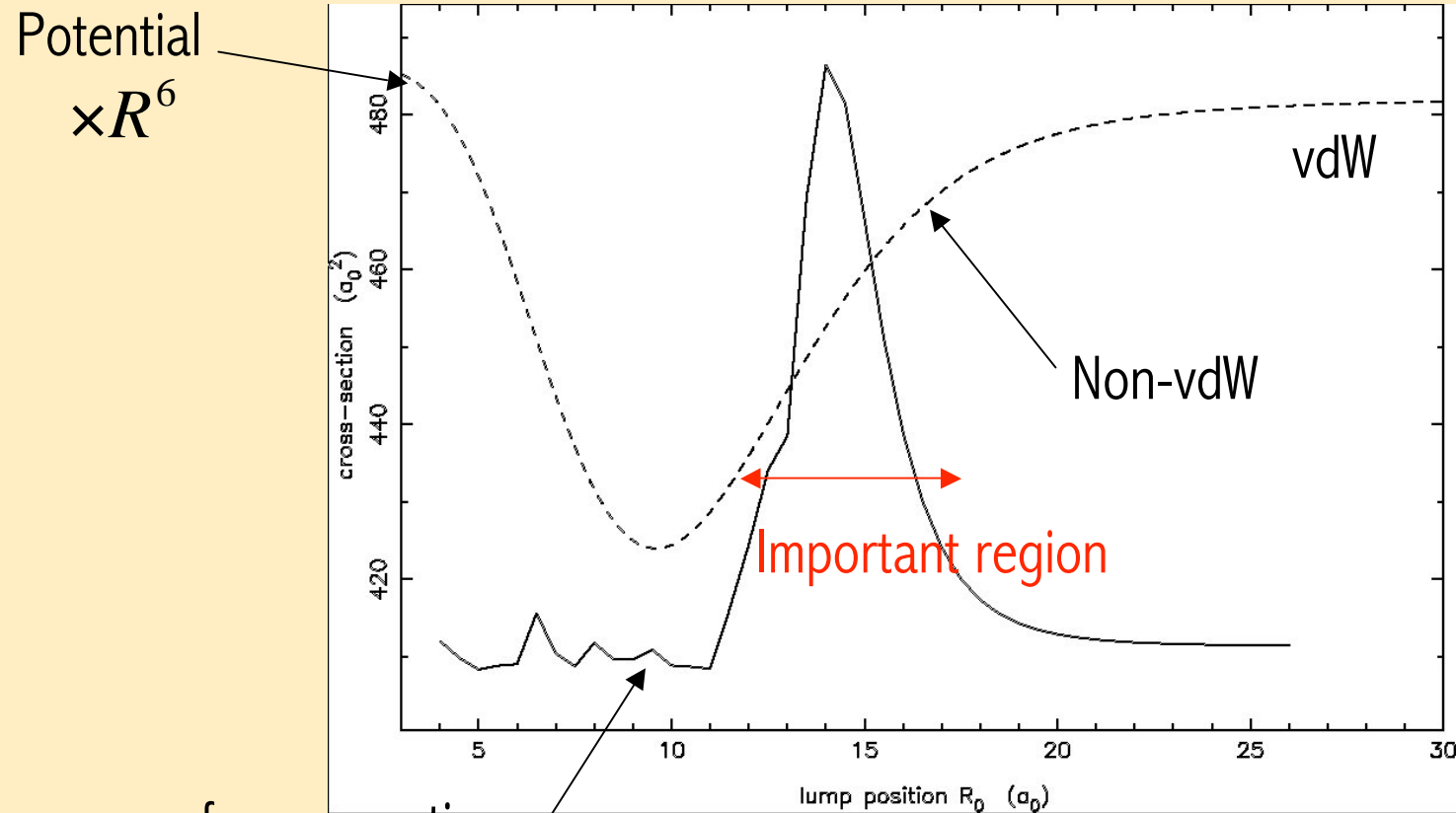
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ABO theory: Why it works

“Intermediate” separations most important:



Response of cross section
To local perturbation in potential

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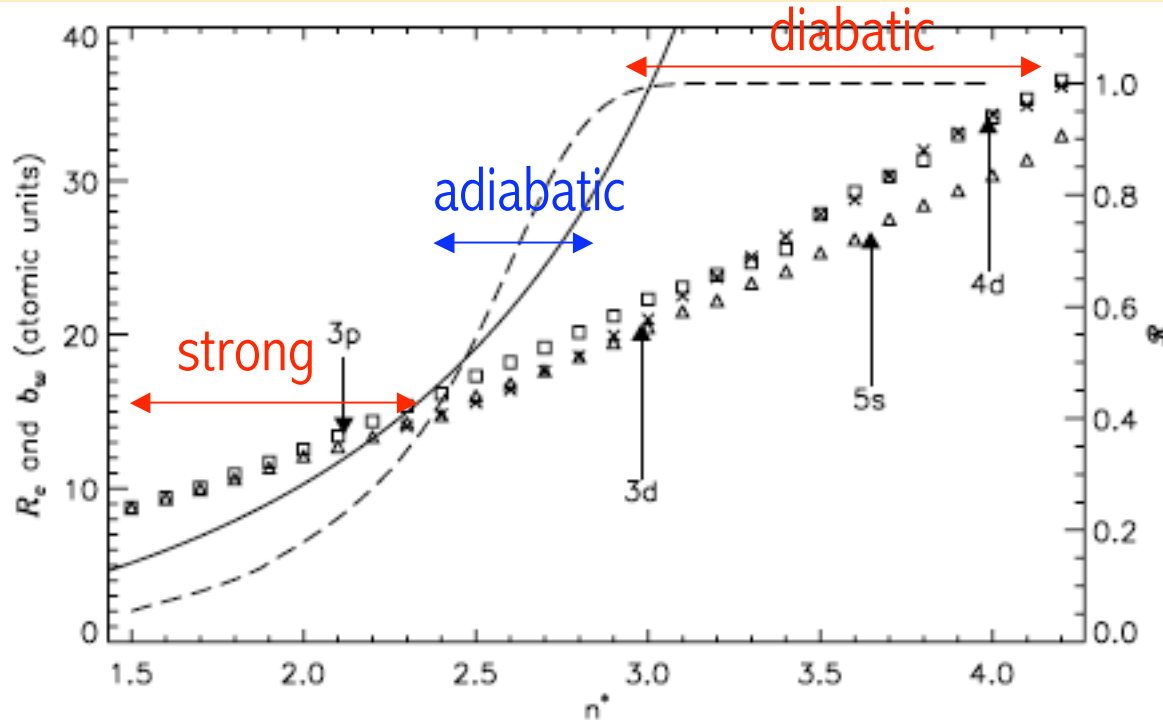


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ABO theory: Why it works

Often, avoided ionic crossings either diabatic, or in the strong collision regime:

Crossing or
Weisskopf
Radius



Impedance to
tunnelling to ionic
configuration

Figure 4. Plots demonstrating various collision regimes. b_w plotted against n^* for s states (triangles), p states (squares) and d states (crosses) for $v = 10^4 \text{ ms}^{-1}$. The full curve is R_e . The reflection coefficient \mathcal{R} , is plotted as the dashed curve.



Summary of Results

- Tables of general data for transitions in neutrals involving s, p, d and f states
- Code for interpolating in tables available
- Table of 4891 strong lines, Li to Ni, mostly neutral + important ionised lines
- Table of 24188 Fe II lines
- Table of 13167 Cr II lines (unpublished)

- Also extended to H Balmer lines



Conclusions

- ABO theory provides a general and widely applicable theory, with accuracy of better than 20%
- Strong lines can now be used with confidence in analysis of cool star spectra (no fudge factors)
- Theory is gaining wide use

- Future: inclusion of ionic state should lead to some improvement

